

A SCALING LIMIT  
WITH MANY  
NONCOMMUTATIVITY  
PARAMETERS

---

1.  $U(N) \rightarrow U(1)^N$  MAGNETIC BACKGROUND
2. CHARGED STRING PROPAGATOR
3. SET OF NONCOMMUTATIVITY  
PARAMETERS
4. STAR PRODUCTS FROM OPE'S
5. SCALING LIMIT

L.D., C. NAPPI      hep-th/0009225

## NEUTRAL STRING IN BACKGROUND B-FIELD

$$\begin{aligned}
 S &= \frac{1}{4\pi\alpha'} \int_{\Sigma} g_{\mu\nu} \partial_{\alpha} X^{\mu} \partial^{\alpha} X^{\nu} - \frac{i}{2} \int_{\Sigma} B_{ij} \epsilon^{\alpha\beta} \partial_{\alpha} X^i \partial_{\beta} X^j \\
 &= \frac{1}{4\pi\alpha'} \int_{\Sigma} g_{\mu\nu} \partial_{\alpha} X^{\mu} \partial^{\alpha} X^{\nu} - \frac{i}{2} \int_{\partial\Sigma} B_{ij} X^i \partial_t X^j
 \end{aligned}$$

HAS SCALING LIMIT = N.C. YANG-MILLS

## CHARGED STRING (WITH A DIFFERENT MAGNETIC FIELD AT EACH END)

$$\begin{aligned}
 S &= \frac{1}{4\pi\alpha'} \int_{\Sigma} g_{\mu\nu} \partial_{\alpha} X^{\mu} \partial^{\alpha} X^{\nu} \\
 &\quad - \frac{i}{2} \int_{-\infty}^{\infty} dt \left( B_{ij}^{(1)} X^i \partial_t X^j \Big|_{\sigma=0} + B_{ij}^{(2)} X^i \partial_t X^j \Big|_{\sigma=\pi} \right)
 \end{aligned}$$

$0 \leq \mu \leq 25$ ,  $D_p$ -branes in  $(0, i)$  directions  
 $1 \leq i \leq p$

SCALING LIMIT = N.C.  $U(1)^N$   
 WITH MASSIVE CHARGED STATES

EQ'S OF MOTION:

$$(\partial_\sigma^2 + \partial_t^2) X^\mu(\sigma, t) = 0$$

 $B_{ij}^{(2)}$  CONSTANT,

 EUCLIDEAN  
 WORLSHEET  
 METRIC
BOUNDARY CONDITIONS:

$$g_{ij} \partial_\sigma X^j + 2\pi i \alpha' B_{ij}^{(1)} X^j \Big|_{\sigma=0} = 0$$

$$g_{ij} \partial_\sigma X^j - 2\pi i \alpha' B_{ij}^{(2)} X^j \Big|_{\sigma=\pi} = 0$$

DIFFER AT EACH END OF THE STRING.

$$\partial_\sigma X^0 \Big|_{\sigma=0, \pi} = 0$$

$$; X^I \Big|_{\sigma=0, \pi} = 0, p+1 \leq I \leq 25.$$

 DIRECTIONS TRANSVERSE  
 TO Dp-branes

 NEUTRAL STRING RESTORED  
 FOR  $B_{ij}^{(1)} = -B_{ij}^{(2)}$ 

 NOW LET  $1 \leq i, j \leq 2.$ 
DZ - DZ SYSTEM

$$B_{12}^{(1)} = q_1 B_{12} \quad ; \quad g_{ij} = g^{-1} \delta_{ij} \quad ; \quad G_{ij} = G^{-1} \delta_{ij}$$

$$B_{12}^{(2)} = q_2 B_{12}$$



# CHARGED STRING NORMAL MODE EXPANSION

---

IN  $X^\pm(\sigma, \tau) \equiv X^\pm(\sigma, \tau)$  BASIS :

---

$$X^+(z, \bar{z}) = x^+ + \frac{i}{2} \sqrt{2\alpha'} \sum_{r \in \mathbb{Z}+A} \frac{a_r}{r} (z^{-r} + \bar{z}^{-r}) - \frac{1}{2} \sqrt{2\alpha'} B \sum_{r \in \mathbb{Z}+A} \frac{a_r}{r} (z^{-r} - \bar{z}^{-r})$$

$$X^-(z, \bar{z}) = x^- + \frac{i}{2} \sqrt{2\alpha'} \sum_{s \in \mathbb{Z}-A} \frac{\tilde{a}_s}{s} (z^{-s} + \bar{z}^{-s}) + \frac{1}{2} \sqrt{2\alpha'} B \sum_{s \in \mathbb{Z}-A} \frac{\tilde{a}_s}{s} (z^{-s} - \bar{z}^{-s})$$

## CHARGED STRING COMMUTATION RELATIONS :

$$[a_r, \tilde{a}_s] = 2\alpha' r \delta_{r, -s} \quad ; \quad [a_r, a_{r'}] = 0$$

$$[\tilde{a}_s, \tilde{a}_{s'}] = 0$$

$$[x^+, x^-] = -\frac{2}{(q_1 + q_2) B_{12}}$$

$$[a_r, x^\pm] = 0$$

$$[\tilde{a}_s, x^\pm] = 0$$

(ACNY)



DEFINITIONS :

$$G = \frac{g}{1+B^2}$$

$$B = g \frac{q_2}{q_1} 2\pi\alpha' B_{12}$$

$$A = \frac{1}{\pi} \left( \arctan B + \arctan \frac{q_2}{q_1} B \right)$$

NON-INTEGGER MODED COMPLEX W.S. BOSONS:

$$a_r : r \in \mathbb{Z} + A$$

$$\tilde{a}_s : s \in \mathbb{Z} - A$$

$$a_r^+ = \tilde{a}_{-r}$$

$$A \rightarrow 0 \quad \sim \quad q_1 \rightarrow -q_2$$

LIMIT BACK TO NEUTRAL STRING :

$$\lim_{A \rightarrow 0} \chi^\pm(z, \bar{z}) = \chi_{n.s.}^\pm(z, \bar{z})$$

$$a_A \rightarrow a_0 \sim p_0$$

$$\text{WHEN } \begin{cases} \lim_{A \rightarrow 0} \left( \chi^+ + i \sqrt{2\alpha'} \frac{a_A}{A} \right) \equiv \chi_0^+ \\ \lim_{A \rightarrow 0} \left( \chi^- - i \sqrt{2\alpha'} \frac{a_{-A}}{A} \right) \equiv \chi_0^- \end{cases}$$

CHARGED STRING OPERATORS :

$$x^\pm, a_A, \tilde{a}_{-A}; a_r, \tilde{a}_s$$

NEUTRAL STRING OPERATORS :

$$x_0^\pm, p_0^+ \equiv \frac{a_0}{\sqrt{2\alpha'}}, p_0^- \equiv \frac{\tilde{a}_0}{\sqrt{2\alpha'}}; a_n, \tilde{a}_m$$

$n, m \in \mathbb{Z}$ .

NEUTRAL STRING COMMUTATION RELATIONS :

$$[a_n, \tilde{a}_m] = 2G n \delta_{n, -m}$$

$$[x_0^+, x_0^-] = 2\Theta^{12}$$

$$[x_0^+, p_0^-] = i2G = [x_0^-, p_0^+]$$

$$[a_n, a_m] = 0$$

$$[\tilde{a}_n, \tilde{a}_m] = 0$$

FOR  $n \neq 0$ :

$$[a_n, x_0^\pm] = 0$$

$$[\tilde{a}_m, x_0^\pm] = 0$$

$$[a_0, x_0^+] = 0$$

$$[\tilde{a}_0, x_0^-] = 0$$

$$\Theta^{12} = -2\pi\alpha' B_G = \frac{-2\pi\alpha'(g)^2 g_{12}}{1 + (g g_{12} 2\pi\alpha' B_{12})^2}$$

## CHARGED STRING SPECTRUM

$$\text{Let } a(z) \equiv \sum_{r \in \mathbb{Z}+A} a_r z^{-r-1}$$

$$\tilde{a}(z) \equiv \sum_{s \in \mathbb{Z}-A} \tilde{a}_s z^{-s-1}$$

VIRASORO ALG.

$$L(z) = \frac{1}{2} G^{-1} : a(z) \tilde{a}(z) : + z^{-2} \frac{1}{2} A(1-A)$$

(CENTRAL CHARGE = 2)

SINCE NO  $\hat{P}^\pm$ ,

LET STATES BE

EIGENSTATES OF POSITION OPERATOR

$$|\chi_+\rangle \equiv e^{-\frac{1}{2}(q_1+q_2)\beta_{12}} \chi_+ \hat{X}_- |0\rangle$$

$$\Rightarrow \hat{X}_+ |\chi_+\rangle = \chi_+ |\chi_+\rangle$$



# NORMAL ORDERING

$$\text{Let } a_r |\chi_+\rangle = 0 \quad r \geq A$$

$$\tilde{a}_s |\chi_+\rangle = 0 \quad s \geq 1-A$$

$$0 \leq A \leq \frac{1}{2}$$

$$L_0 |\chi_+\rangle = \frac{1}{2} A(1-A) |\chi_+\rangle$$

FOR EACH OF THE INFINITE  
NUMBER OF STATES  $|\chi_+\rangle$ .

## OSCILLATOR TOWER :

STATE

$$\alpha' M^2 = L_0 - 1$$

$|\chi_+\rangle$

$a_{-1+A} |\chi_+\rangle$

$\tilde{a}_{-1-A} |\chi_+\rangle$

⋮

$$\begin{aligned} & -1 + \frac{1}{2} A(1-A) < 0 \\ & \left[ \begin{aligned} & -1 + \frac{1}{2} A(1-A) + (1-A) \\ & = -\frac{1}{2} A(1+A) < 0 \end{aligned} \right. \\ & \left[ \begin{aligned} & -1 + \frac{1}{2} A(1-A) + (1+A) \\ & = \frac{1}{2} A(3-A) > 0 \end{aligned} \right. \end{aligned}$$

POLARIZATIONS  
ARE SPLIT IN FREQUENCY:

- 1 TACHYONIC MODE
- 1 MASSIVE MODE

LANDAU LEVELS:

$$(a_A |x_+\rangle = 0)$$

$$|x_+\rangle, \hat{a}_{-A} |x_+\rangle, \hat{a}_{-A} \hat{a}_{-A} |x_+\rangle, \dots$$

$$\Delta(L_0 - 1) = A$$

SO THESE STATES HAVE FREQUENCY SEPARATION A.

THEY ARE LIKE EQUALLY SPACED LANDAU LEVELS (EACH) OF INFINITE DEGENERACY.

AND EACH IS THE LOWEST STATE OF AN OSCILLATOR TOWER.

TO COMPUTE  
CHARGED STRING PROPAGATOR.

CONSIDER THE COHERENT STATES:

$$|\alpha\rangle = e^{-\frac{1}{2}(\alpha_1 + \alpha_2)\beta_{12}} x^- |0\rangle$$

$$\langle\beta| = \langle 0| e^{-\frac{1}{2}(\alpha^+ + i\frac{\sqrt{2}\alpha^+}{\alpha} a_A)}$$

THIS WILL ENSURE  
CHARGED PROPAGATOR  $\rightarrow$  NEUTRAL PROP.  
 $A \rightarrow 0$

$$\text{SINCE } \lim_{A \rightarrow 0} \langle\beta| \chi_0^+ \chi_0^- + \chi_0^- \chi_0^+ |\alpha\rangle = 0.$$

$$\oint N \chi_0^+ \chi_0 \equiv \frac{1}{2}(\chi_0^+ \chi_0^- + \chi_0^- \chi_0^+)$$

IN NEUTRAL STRING.



## CHARGED STRING PROPAGATOR:

For  $|z\rangle|\bar{z}\rangle$ :

$$\begin{aligned}
 \langle \chi^+(z, \bar{z}) | \chi^-(\zeta, \bar{\zeta}) \rangle &\equiv \langle \beta | \chi^+(z, \bar{z}) \chi^-(\zeta, \bar{\zeta}) | \alpha \rangle \\
 &= \frac{-2\alpha' \pi g}{\beta + \frac{q_2}{q_1} \beta} - 2\alpha' G \frac{1}{A} (\zeta^A + \bar{\zeta}^A - 1) + 2i\alpha' G \beta \frac{1}{A} (\zeta^A - \bar{\zeta}^A) \\
 &\quad + \alpha' G \left[ f\left(\frac{\zeta}{2}\right) + f\left(\frac{\bar{\zeta}}{2}\right) + f\left(\frac{\zeta}{2}\right) + f\left(\frac{\bar{\zeta}}{2}\right) \right] \\
 &\quad + \alpha' G \beta^2 \left[ f\left(\frac{\zeta}{2}\right) + f\left(\frac{\bar{\zeta}}{2}\right) - f\left(\frac{\zeta}{2}\right) - f\left(\frac{\bar{\zeta}}{2}\right) \right] \\
 &\quad + 2i\alpha' G \beta \left[ -f\left(\frac{\zeta}{2}\right) + f\left(\frac{\bar{\zeta}}{2}\right) \right]
 \end{aligned}$$

where

$$f(\rho) \equiv \sum_{\substack{r=n+A \\ n \geq 0}} \frac{\rho^r}{r} \quad ; \quad \lim_{A \rightarrow 0} f(\rho) = -\ln(1-\rho) + \lim_{A \rightarrow 0} \frac{\Delta^A}{A}$$

(INCOMPLETE BETA FUNCTION)

OTHER NON ZERO COMPONENT:  $|z\rangle\langle s|$

$$\langle \chi^-(z, \bar{z}) \chi^+(s, \bar{s}) \rangle \equiv \langle \beta | \chi^-(z, \bar{z}) \chi^+(s, \bar{s}) | \alpha \rangle$$

$$= \frac{2\alpha' \pi g}{B + \frac{q_2}{q_1} B} - 2\alpha' G \frac{1}{A} (z^A + \bar{z}^A - 1) + 2i\alpha' G \frac{B}{A} (z^A - \bar{z}^A)$$

$$+ \alpha' G \left[ g\left(\frac{s}{z}\right) + g\left(\frac{\bar{s}}{\bar{z}}\right) + g\left(\frac{s}{\bar{z}}\right) + g\left(\frac{\bar{s}}{z}\right) \right]$$

$$+ \alpha' G B^2 \left[ g\left(\frac{s}{z}\right) + g\left(\frac{\bar{s}}{\bar{z}}\right) - g\left(\frac{s}{\bar{z}}\right) - g\left(\frac{\bar{s}}{z}\right) \right]$$

$$- 2i\alpha' G B \left[ -g\left(\frac{s}{z}\right) + g\left(\frac{\bar{s}}{\bar{z}}\right) \right].$$

$$g(\rho) \equiv \sum_{s=n-A}^{\rho^s} \frac{\rho^s}{s} \quad ; \quad \lim_{A \rightarrow 0} g(\rho) = -\ln(1-\rho).$$

COMPARISON WITH  
NEUTRAL STRING PROPAGATOR:

$$\begin{aligned}
 & \langle X^i(z, \bar{z}) X^j(\zeta, \bar{\zeta}) \rangle_{\text{N.S.}} \\
 &= -\alpha' \left[ \frac{1}{2} g^{ij} \ln(z - \zeta) + \frac{1}{2} g^{ij} \ln(\bar{z} - \bar{\zeta}) \right. \\
 & \quad \left. + \left( -\frac{1}{2} g^{ij} + G^{ij} + \frac{\Theta^{ij}}{2\pi\alpha'} \right) \ln(z - \bar{\zeta}) \right. \\
 & \quad \left. + \left( -\frac{1}{2} g^{ij} + G^{ij} - \frac{\Theta^{ij}}{2\pi\alpha'} \right) \ln(\bar{z} - \zeta) \right. \\
 & \quad \left. - \frac{i}{2\alpha'} \Theta^{ij} \right].
 \end{aligned}$$

CHARGED PROPAGATOR FOR ALL  $z, \zeta$ :

$$\begin{aligned}
 G^{+-}(z, \bar{z}; \zeta, \bar{\zeta}) &= \Theta(|z| - |\zeta|) \langle X^+(z, \bar{z}) X^-(\zeta, \bar{\zeta}) \rangle \\
 & \quad + \Theta(|\zeta| - |z|) \langle X^-(\zeta, \bar{\zeta}) X^+(z, \bar{z}) \rangle
 \end{aligned}$$

$$\text{AND } G^{-+}(z, \bar{z}; \zeta, \bar{\zeta}) = G^{+-}(\zeta, \bar{\zeta}; z, \bar{z}).$$

$$\Rightarrow \lim_{A \rightarrow 0} G^{+-} = \langle X^+ X^- \rangle_{\text{N.S.}}$$

$$\lim_{A \rightarrow 0} G^{-+} = \langle X^- X^+ \rangle_{\text{N.S.}}$$



## COMPUTATION OF NONCOMMUTATIVITY PARAMETERS :

1. CALCULATE THE EQUAL TIME COMMUTATOR VIA A SHORT DISTANCE EXPANSION OF THE PROPAGATOR :
2. EVALUATE PROPAGATOR ON EACH BOUNDARY

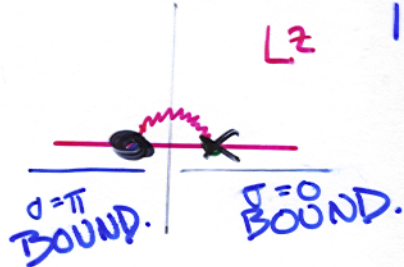
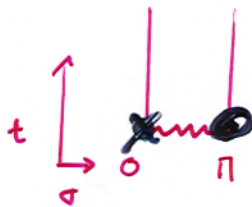
$$\Rightarrow 2\Theta^{12} \equiv [x^+(\tau), x^-(\tau)]$$

$$\equiv T (x^+(\tau) x^-(\tau^-) - x^+(\tau) x^-(\tau^+))$$

$$= \lim_{\epsilon \rightarrow 0} \langle x^+(\tau) x^-(\tau - \epsilon) \rangle - \langle x^-(\tau + \epsilon) x^+(\tau) \rangle$$

↑  
PROPAGATOR S

FOR EXAMPLE :



ON  $\sigma = 0$  BOUNDARY

$$\Rightarrow z = |z| = \rho > 0, \quad s = |s| = \rho' > 0$$

ON THE BOUNDARY  $\sigma = 0$   
THE PROPAGATOR IS :

$$\begin{aligned} \langle X^+(z, \bar{z}) X^-(s, \bar{s}) \rangle \Big|_{\sigma=0} &= -2 \frac{\alpha' \pi g}{B + \frac{q_2}{q_1} B} \\ &+ 4\alpha' G \sum_{n=0}^{\infty} \frac{1}{n+A} \left(\frac{s}{z}\right)^{n+A} \\ &+ \frac{2\alpha' G}{A} - \frac{4\alpha' G}{A} s^A \end{aligned}$$

$$\begin{aligned} \langle X^-(z, \bar{z}) X^+(s, \bar{s}) \rangle \Big|_{\sigma=0} &= \frac{2\alpha' \pi g}{B + \frac{q_2}{q_1} B} \\ &+ 4\alpha' G \sum_{n=1}^{\infty} \frac{1}{n-A} \left(\frac{s}{z}\right)^{n-A} \\ &+ \frac{2\alpha' G}{A} - \frac{4\alpha' G}{A} z^A \end{aligned}$$

$$\begin{aligned}
 & [X^+(\tau), X^-(\tau)] \Big|_{\tau=0} \\
 &= \lim_{\epsilon \rightarrow 0} \left( \frac{-4\alpha' \pi g}{\beta + \frac{g_2}{g_1} \beta} + 4\alpha' G \left[ \sum_{n=0}^{\infty} \frac{1}{n+A} \left( \frac{\tau-\epsilon}{\tau} \right)^{n+A} - \sum_{n=1}^{\infty} \frac{1}{n-A} \left( \frac{\tau}{\tau+\epsilon} \right)^{n-A} \right] \right) \\
 &= \frac{-4\alpha' \pi g}{\beta + \frac{g_2}{g_1} \beta} + 4\alpha' G \pi \cot \pi A \\
 &= \boxed{\frac{-4\alpha' \pi (g)^2}{q_1} \frac{2\pi \alpha' \beta_{12}}{1 + (g q_1 2\pi \alpha' \beta_{12})^2}} \\
 &\equiv 2\ominus^{12}
 \end{aligned}$$

THE NON COMMUTATIVITY  
PARAMETER AT  $\tau=0$



AT THE  $\sigma = \pi$  BOUNDARY,

$$z = |z| e^{i\pi} = \tau < 0, \quad \bar{z} = |\bar{z}| e^{i\pi} = \tau' < 0$$

$$\Rightarrow [\chi^+(\tau), \chi^-(\tau)] \Big|_{\sigma = \pi}$$

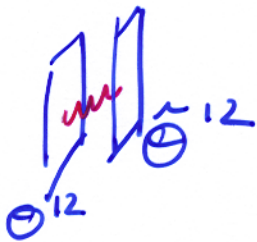
$$= \lim_{\epsilon \rightarrow 0} \left( \langle \chi^+(\tau) \chi^-(\tau + \epsilon) \rangle - \langle \chi^-(\tau - \epsilon) \chi^+(\tau) \rangle \right) \Big|_{\sigma = \pi}$$

$$= \frac{-4\alpha' \pi (g)^2 g_2 \pi \alpha' B_{12}}{1 + (g)^2 (g_2 2\pi \alpha' B_{12})^2}$$

$$\approx 2\tilde{\Theta}^{12}$$

THE NON COMMUTATIVITY PARAMETER  
AT  $\sigma = \pi$  BOUNDARY.

FOR COINCIDENT D2 BRANES,

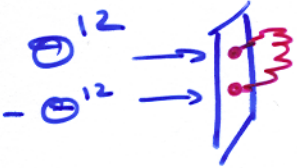


EACH BRANE HAS ITS  
OWN  
NON COMMUTATIVITY  
PARAMETER.

FOR NEUTRAL STRING LIMIT:

$$q_1 = -q_2$$

$$\Rightarrow \tilde{\Theta}^{12} = -\Theta^{12}$$



BOTH ENDS OF THE STRING  
ARE ON THE SAME D2-BRANE.

THE NC PARAMETER AT ONE END  
OF THE STRING IS EQUAL TO  
MINUS THAT AT THE OTHER END.

# SHORT DISTANCE BEHAVIOR AND STAR PRODUCTS :

OPE FOR TACHYON VERTEX OP.'S :  
(ON  $\sigma = 0$  BOUNDARY)

$$\begin{aligned}
 & e^{ip \cdot X(\tau)} e^{iq \cdot X(\tau')} \\
 &= (\tau - \tau')^{4\alpha' G (p \cdot q_+ + p \cdot q_-)} \\
 & \cdot e^{-\Theta'^2 (p \cdot q_+ - p \cdot q_-)} : e^{i(p+q) \cdot X(\tau')} :
 \end{aligned}$$

IN SCALING LIMIT ( $\alpha' \rightarrow 0$ ) ( $G, \Theta'^2$  FIXED)

$$e^{ip \cdot X(\tau)} e^{iq \cdot X(\tau')} \sim e^{ip \cdot X(\tau')} * e^{iq \cdot X(\tau')}$$

STAR PRODUCT :

$$\begin{aligned}
 f(x) * g(x) &= e^{\frac{i}{2} \Theta^{ij} \frac{\partial}{\partial x^i} \frac{\partial}{\partial x^j}} f(x+\xi) g(x+\xi) \Big|_{\xi=0} \\
 &= fg + \frac{i}{2} \Theta^{ij} \partial_i f \partial_j g + o(\Theta^2)
 \end{aligned}$$

⇒ A DIFFERENT STAR PRODUCT FOR EACH BOUNDARY.



THE SPECTRUM IN THE SCALING LIMIT:

$$g^{-1} \rightarrow \epsilon, \quad \alpha' \rightarrow \sqrt{\epsilon} \quad \text{AS } \epsilon \rightarrow 0$$

$$(\alpha' \beta_{12} \rightarrow \sqrt{\epsilon} \text{ ; keeping } \beta_{12} \text{ fixed}).$$

$$\Rightarrow \Theta^{12} \rightarrow \frac{1}{q_1 \beta_{12}}, \quad \hat{\Theta}^{12} \rightarrow \frac{1}{q_2 \beta_{12}}, \quad G \rightarrow \frac{1}{(q_1 2\pi)^2}.$$

THESE PARAMETERS ARE FINITE  
IN THE SCALING LIMIT.

$$A \rightarrow - \frac{(q_1 + q_2) \sqrt{\epsilon}}{2\pi^2 q_1 q_2}$$

MASSES OF  
CHARGED VECTOR STATES :

$$\begin{aligned} \alpha_{-1+A} |X_+ \rangle \quad m^2 &= -\frac{1}{2\alpha'} A(1+A) \\ &\rightarrow \frac{1}{2} \frac{(q_1 + q_2) \beta_{12}}{2\pi^2 q_1 q_2} = \frac{(q_1 + q_2) G}{q_2 \Theta^{12}} \end{aligned}$$

$$\begin{aligned} \hat{\alpha}_{-1-A} |X_+ \rangle \quad m^2 &= \frac{1}{2\alpha'} A(3-A) \\ &\rightarrow -\frac{3}{2} \frac{(q_1 + q_2) \beta_{12}}{2\pi^2 q_1 q_2} = -\frac{3(q_1 + q_2) G}{q_2 \Theta^{12}}. \end{aligned}$$

CHARGED VECTOR STATES THAT SURVIVE IN THE SCALING LIMIT:

$$\begin{aligned}
 & a_{-1+A} |X_+\rangle, \tilde{a}_{-1-A} |X_+\rangle \\
 & \downarrow \\
 & a_{-1+A} \hat{a}_{-A} |X_+\rangle, \tilde{a}_{-1-A} \hat{a}_{-A} |X_+\rangle \\
 & \vdots \\
 & \text{VECTORS AT ALL LANDAU LEVELS.}
 \end{aligned}$$

COMPLETE SPECTRUM OF THE  $U(1)^N$  NON COMMUTATIVE GAUGE THEORY

( $U(N)$  GAUGE THEORY EXPANDED AROUND A  $U(1)^N$  BACKGROUND) :

$N$  NEUTRAL MASSLESS GLUONS  
 $N^2 - N$  CHARGED VECTOR SECTORS

