

# Non-Commutative Geometry & Yukawa Couplings in F-Theory

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29/10/09 UNC

# Outline

- Motivation
- F-Theory
- Yukawa in Commutative Geometry
- More F-Theory
- Yukawa in Non-Commutative Geometry
- Conclusions and Applications

# Motivation

- Structures of Yukawa's in F-theory: what types are allowed?
- Standard Model/MSSM  $\subset$  String Theory?  
Flavor Hierarchy from??

# In This Talk

- Yakawa Coupling in 7-brane Gauge Theory
- Possible Relevance for Hierarchy in Quark Sector

# Flavor Hierarchy

$$W_{MSSM} = m_u^{ij} Q^i \bar{U}^j H_u + m_d^{ij} Q^i \bar{D}^j H_d + \dots$$

$i, j$ : Flavor indices

**Diagonalize:**  $V_L \cdot m \cdot V_R^\dagger = \text{diag}(m_1, m_2, m_3) \sim \text{diag}(m_1, 0, 0)$

**CKM Matrix:**  $|V_{CKM}| = V_{L,u} \cdot V_{L,d}^\dagger \sim \mathbf{I}$

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hierarchical !

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we explain WHY naturally

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$m^{ij}$  : rank 1 + corrections



# F-theory

Structure          Dimensions

gravity                  10

gauge fields (7-brane)      8

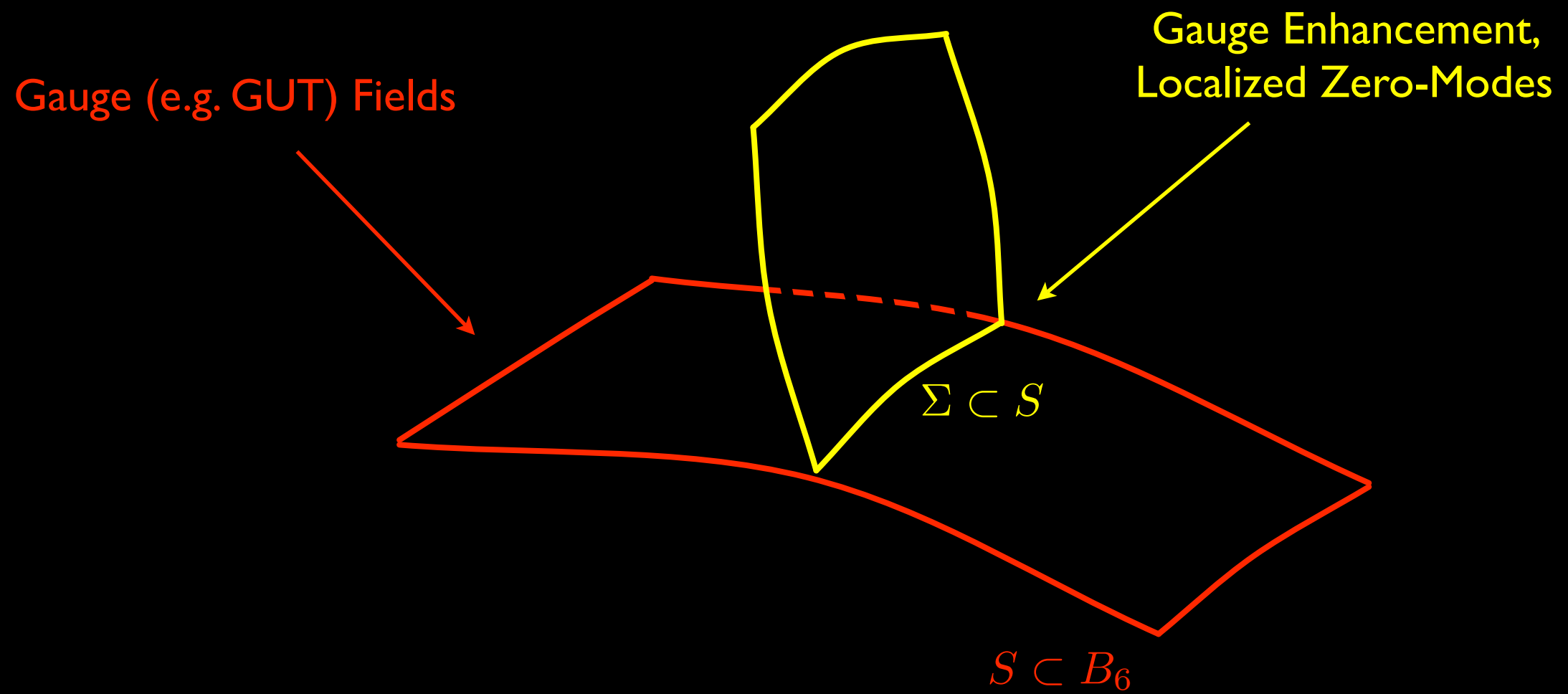
matter fields ( $7 \cap 7'$ )      6

Yuk. couplings ( $7 \cap 7' \cap 7''$ )      4

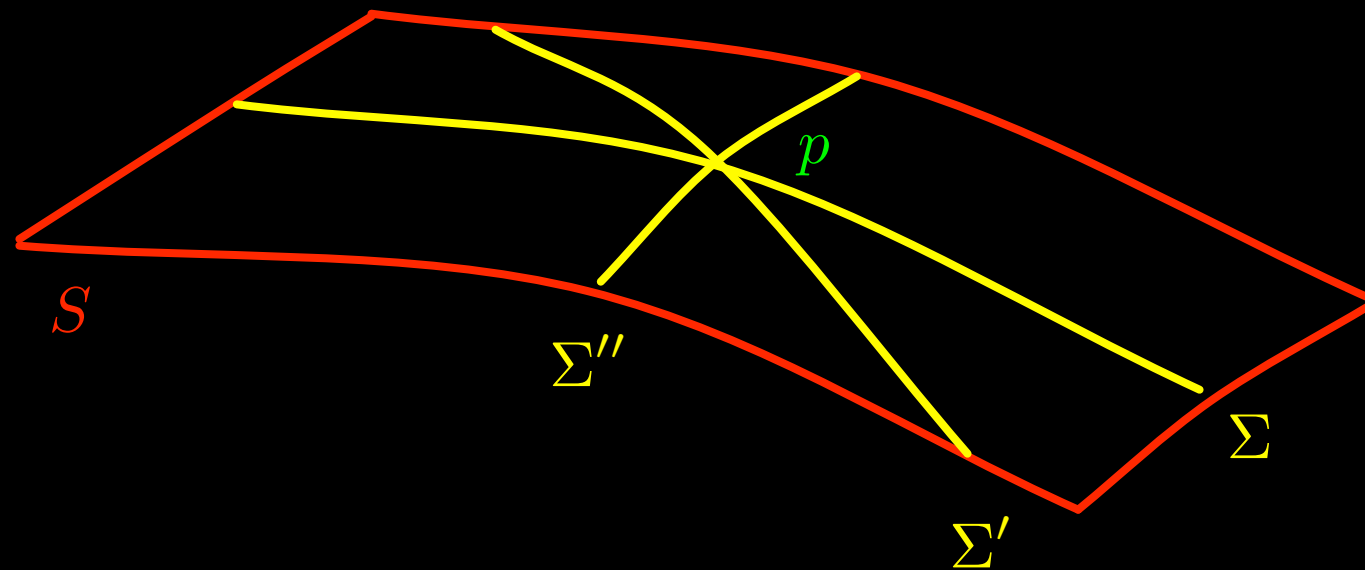
# F-theory

Structure	Dimensions	Dimensions in $B_6 \subset CY_4$
gravity	10	6
gauge fields (7-brane)	8	4 ( $S$ )
matter fields ( $7 \cap 7'$ )	6	2 ( $\Sigma$ )
Yuk. couplings ( $7 \cap 7' \cap 7''$ )	4	0 ( $p$ )

# F-theory



# F-theory

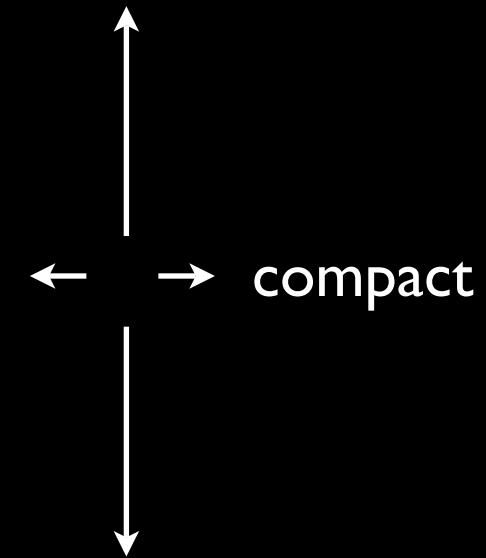


$$G_p \supset G_\Sigma \supset G_S$$

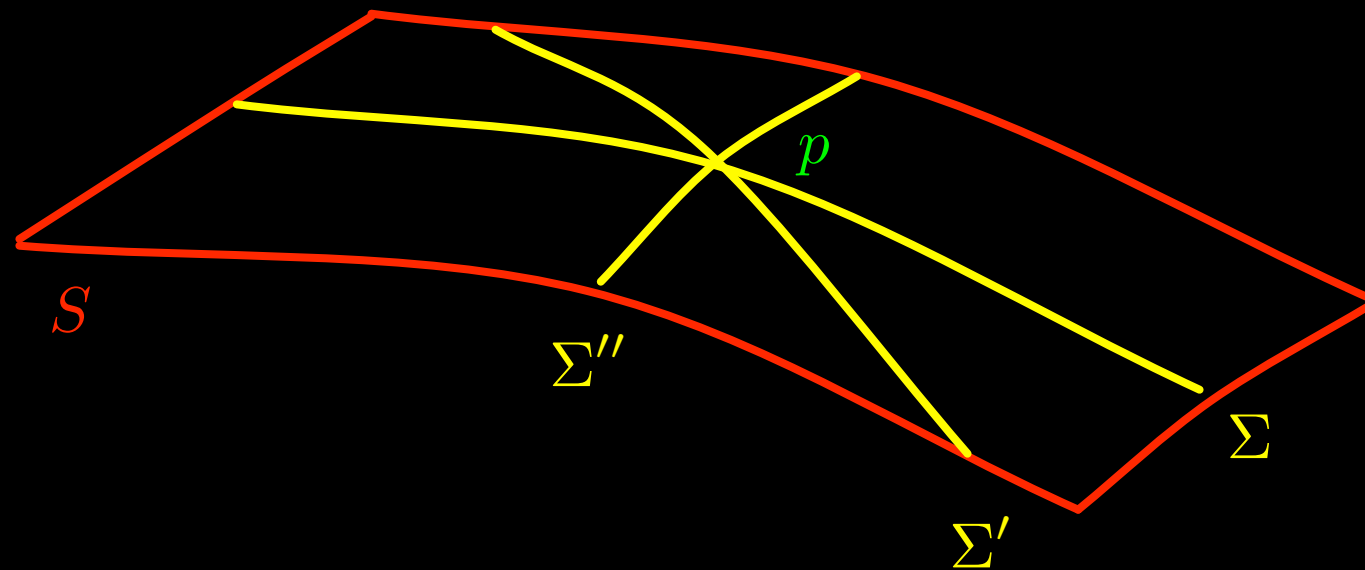
# F-theory

**Decoupling  
Assumption:**

non-compact



$$\frac{M_{GUT}}{M_{Pl}} \sim 10^{-3} \ll 1$$

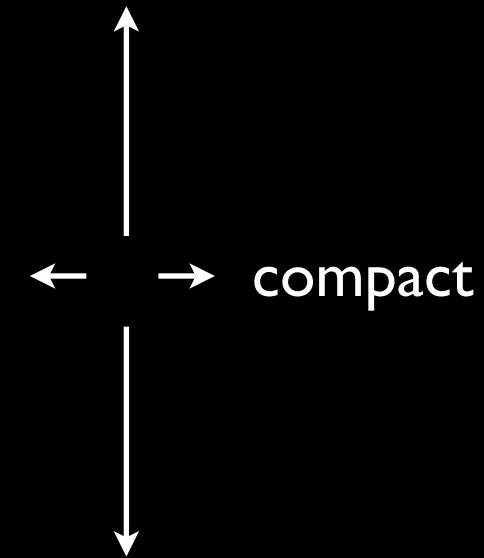


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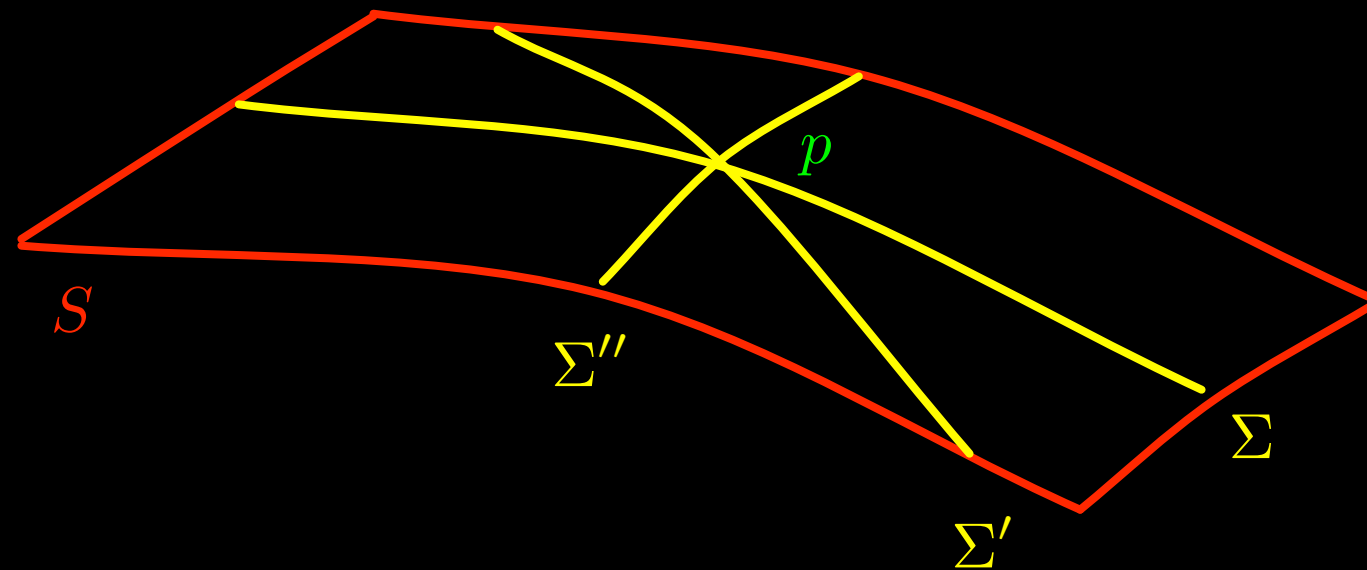
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$\Rightarrow$  Yukawa coupling described by gauge theory living on  $S$  of gauge group  $G_p$  Higgsed down to  $G_S$ .

# III. Commutative 7-Brane Theory

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$\Rightarrow$  Rank 1 Yukawa



# Superpotential

$$W = \int_S \text{Tr} (\varphi \wedge F^{(0,2)})$$

[Beasley-Heckman-Vafa, Donagi-Wijnholt]

$$F^{(0,2)} = \bar{\partial}\bar{A} + \bar{A} \wedge \bar{A} \quad , \quad \bar{A} = A^{(0,1)}$$

$$\varphi^{(2,0)} : \text{adj}(G)$$

cf. D9 in pert. IIB  $\rightarrow$  hol. C-S

[Witten]

$$W = \int_{CY_3} \text{Tr} \left( \Omega \wedge \left( \bar{A}\bar{\partial}\bar{A} + \frac{2}{3}\bar{A} \wedge \bar{A} \wedge \bar{A} \right) \right)$$

dim. reduce  $\quad \Omega\bar{A}_\perp \rightarrow \varphi^{(0,2)}$

## F-term Eqn

$$F^{(0,2)} = \bar{\partial}\bar{A} + \bar{A} \wedge \bar{A} = 0$$

$$\bar{\partial}_A \varphi = 0$$

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## ~~D-term Eqn~~

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$W$  independent of  $g_{i\bar{j}}$ ,  $A^{(1,0)}$ ,  $\bar{\varphi}$

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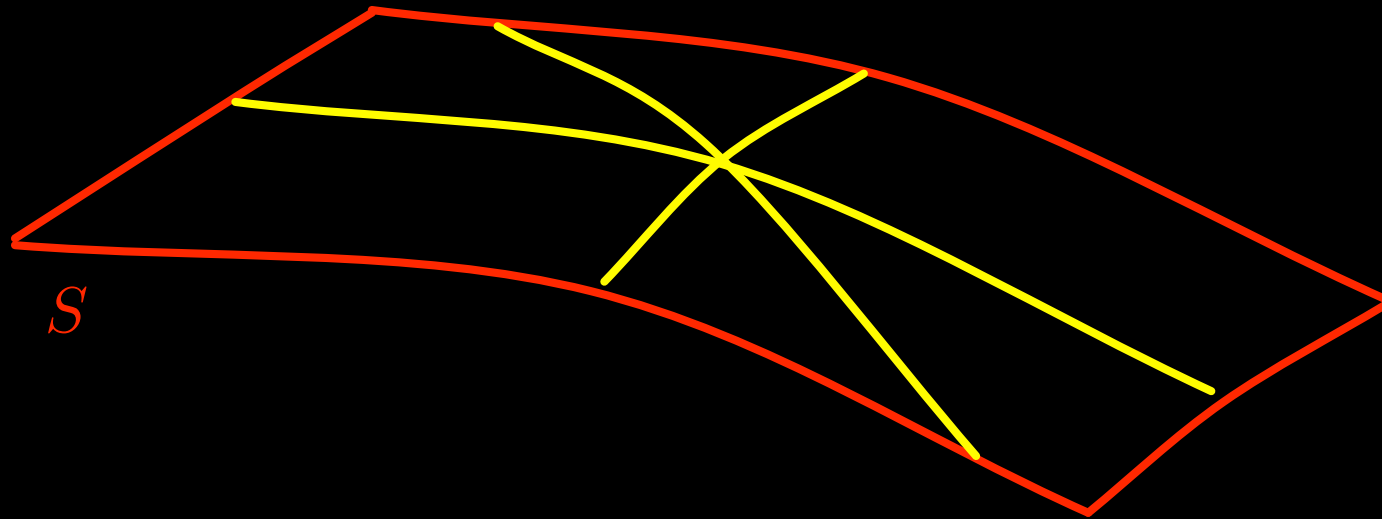
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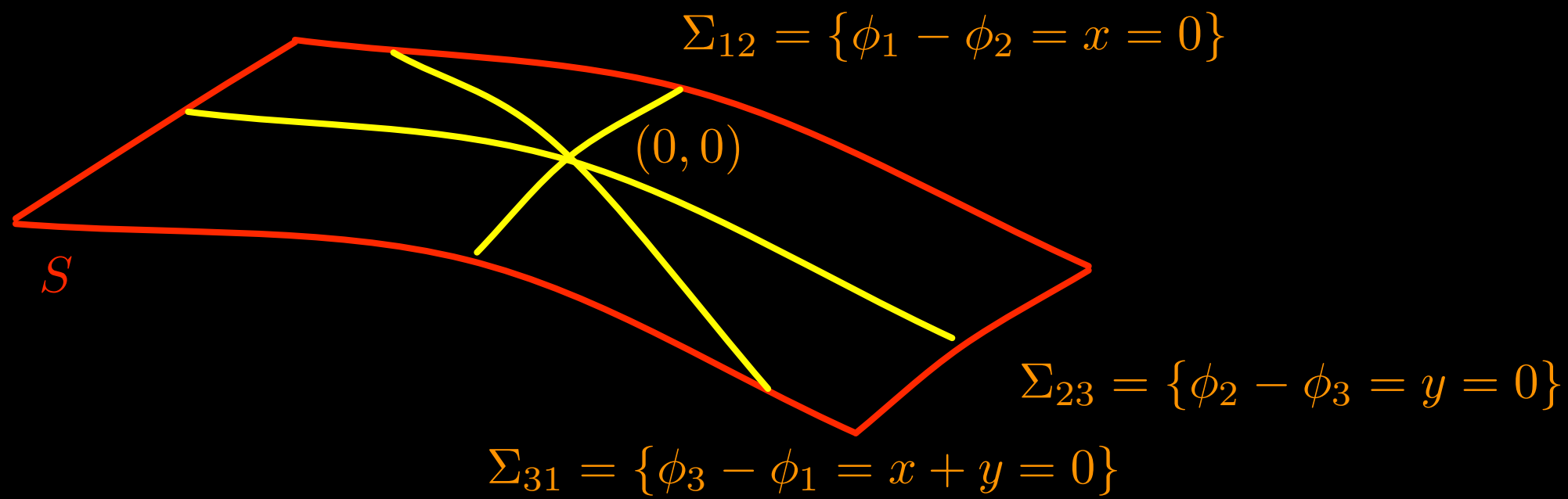
“deformation” away from the given background  $\Rightarrow$  chiral zero-modes localized on matter curves

$$W = \int \text{Tr} \left( \varphi \wedge F^{(0,2)} \right) = \int \text{Tr} \left( \varphi^{(1)} \wedge \bar{A}^{(1)} \wedge \bar{A}^{(1)} \right) \neq 0$$

$\Leftrightarrow$  obstruction to deformation



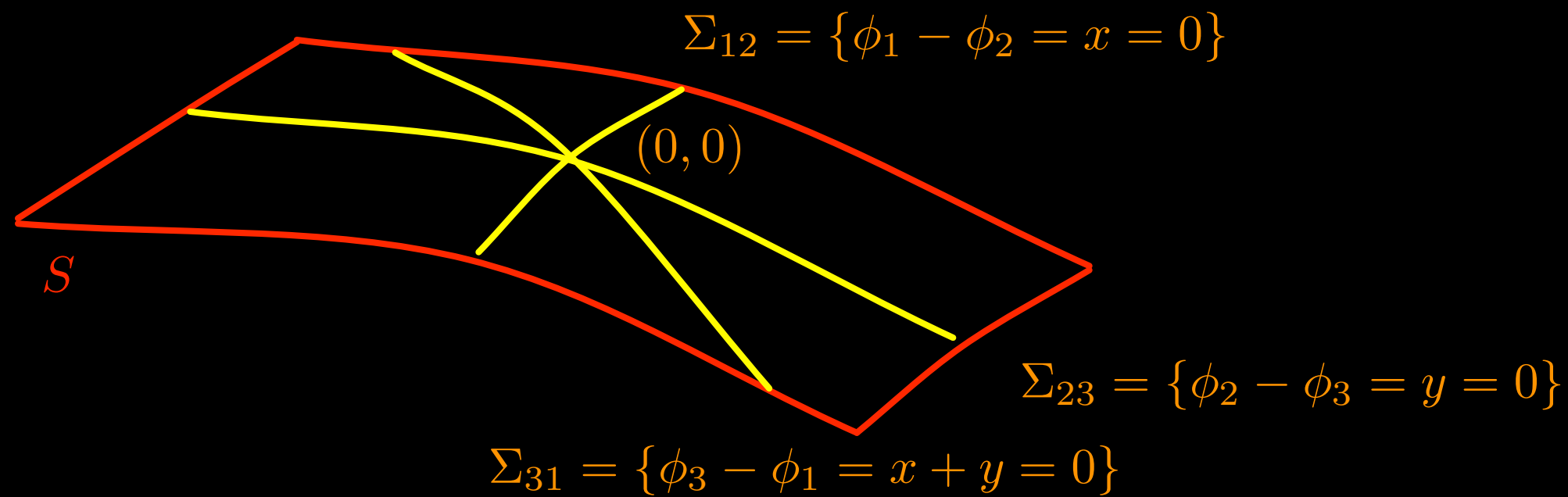
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eg.  $G = U(3)$

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**Note: turning on background  $F^{(1,1)}$  won't change  $W$ !**

F-term zero-mode eqn  $\Rightarrow$

$$\bar{A}_{ab}^{(1)} = \bar{\partial} \left( \frac{\phi_{ab}^{(1)} - h_{ab}}{\phi_a - \phi_b} \right), \quad \phi_{ab}^{(1)} \Big|_{\Sigma_{ab}} = h_{ab} \Big|_{\Sigma_{ab}} \quad \text{hol.}$$

Gauge-inequivalent zero-modes given by

$$\mathcal{O} / \{\phi_a - \phi_b\}, \quad \text{namely} \quad h_{ab} \sim h_{ab} + \tilde{h}_{ab} (\phi_a - \phi_b)$$

captured by a cohomology defined for given background!

Yukawa = pairing between cohomology classes

# Yukawa as Residues

for example, take  $h_{31} = 1$  (Higgs field)

$$\begin{aligned}
 W^{(ij)} &= \oint dx \oint dy \frac{h_{12}^{(i)} h_{23}^{(j)}}{x y} \\
 &= 0 \quad \text{unless } h_{12}^{(i)}, h_{23}^{(j)} \text{ both constants}
 \end{aligned}$$

$\Rightarrow$  Rank 1

in general Rank =  $\#(\Sigma_{12} \cap \Sigma_{23})$

**NEED BETTER!!**

Need to consider  $H = H_R + \tau H_{NS}$

$$\int_{B_6} H_R \wedge H_{NS} + N_{D3} = \frac{1}{24} \chi(CY_4)$$



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$\Rightarrow$  Non-Commutativity in Geometry

$$\begin{aligned} [X^i, X^j] &= \theta^{ij}(X) \longleftarrow \text{holomorphic, Poisson} \\ \nabla_k \theta^{ij} &= H_{ij}^k \end{aligned}$$

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(pert. IIB, flat space, R-flux)

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Myer's Effect

$$V_{D0}^{flux} = \text{Tr}(X^I X^J X^K) G_{0IJK} + \dots$$

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T-dual

$$V_{D3}^{flux} = \text{Tr}(X^j X^{\bar{j}} X^{\bar{k}}) H_{i\bar{j}\bar{k}} + \dots$$

$$\Rightarrow \nabla_k \theta^{ij} = H^{ij}_k + \dots$$

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deforming B-model by elements of  $H^0(\Lambda^2 T^{1,0})$  [Kapustin, K-Li]

$$\Omega \rightarrow \Omega + \theta' \quad , \quad \theta'_k = \theta^{ij} \Omega_{ijk}$$

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open B-model amplitude  $\Rightarrow$  N-C with  $\theta$  [Pestun, Cattaneo-Felder]

relation between generalized Kähler geometry and world-sheet B-field fluxes [Gates, Hull, Rocek]

$$\Rightarrow \nabla \theta = H + \dots$$

(pert. IIB)

## Bulk H-fluxes

⇒ N-C deformation of 7-brane Gauge Theory



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Q: How does Yuk get deformed?

# non-commutativity of surface $S$

$$[x, y] = \theta(x, y) \quad \text{on local patch } \mathbb{C}^2$$

constant  $\theta \rightarrow$  Moyal  $\star$  product

$$\begin{aligned} f \star g &= e^{\theta^{ij} \partial_i \wedge \partial_j} f \otimes g \\ &= fg + \theta^{ij} \partial_i f \partial_j g + \mathcal{O}(\theta^2) \end{aligned}$$

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covariantize to  $\circledast$  for general  $\theta(x, y)$  and  
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 for functions and forms

pert. series in  $\theta$

deformed Superpotential:  $\wedge \rightarrow \circledast$

$$W = \int_S \text{Tr} \left( \varphi \circledast (\bar{\partial} \bar{A} + \bar{A} \circledast \bar{A}) \right)$$

complex gauge transformation

$$\bar{A} \rightarrow g^{-1 \circledast} \circledast \bar{A} \circledast g + g^{-1 \circledast} \circledast \bar{\partial} g \quad , \quad \varphi \rightarrow g^{-1 \circledast} \circledast \varphi \circledast g$$

$$g^{-1 \circledast} \circledast g = g \circledast g^{-1 \circledast} = 1 \quad , \quad g^{-1 \circledast} = g^{-1} + \mathcal{O}(\theta)$$

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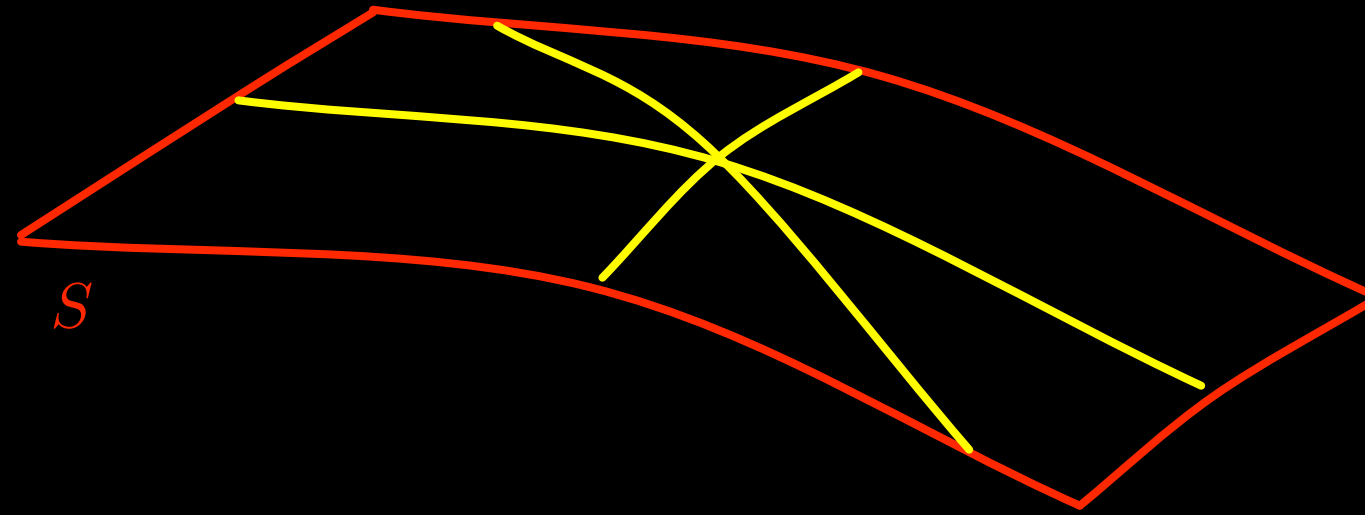
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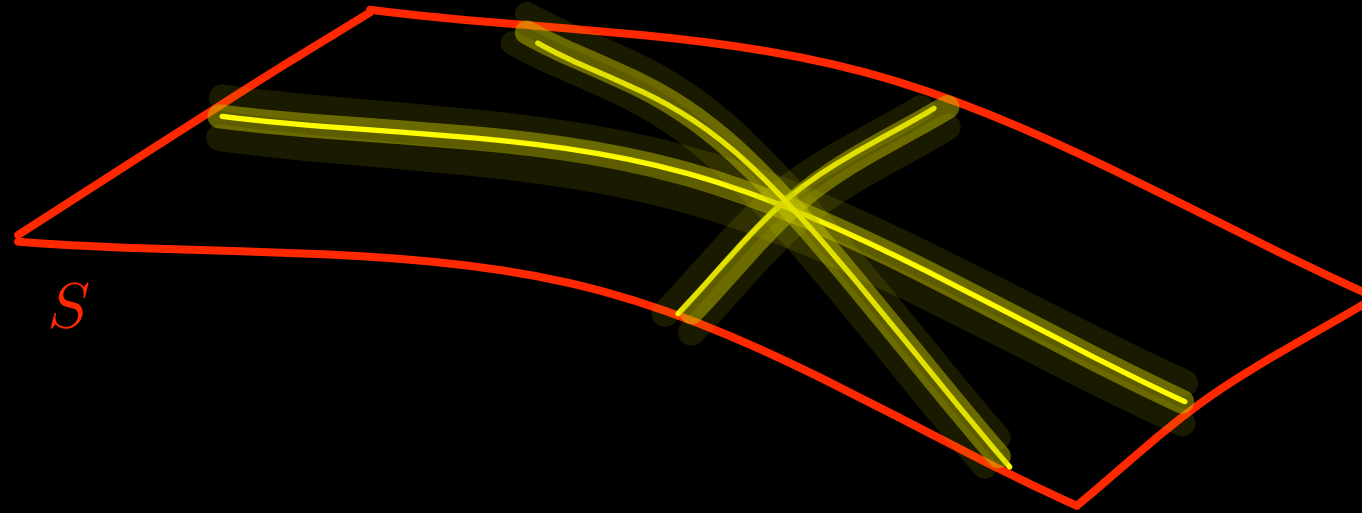
$$g^{-1 \circledast} \circledast g = g \circledast g^{-1 \circledast} = 1 \quad , \quad g^{-1 \circledast} = g^{-1} + \mathcal{O}(\theta)$$

$\bar{A}, \varphi$  take value now in Universal Enveloping Algebra

Matter Curves  $\Rightarrow$  Fuzzy Matter Curves

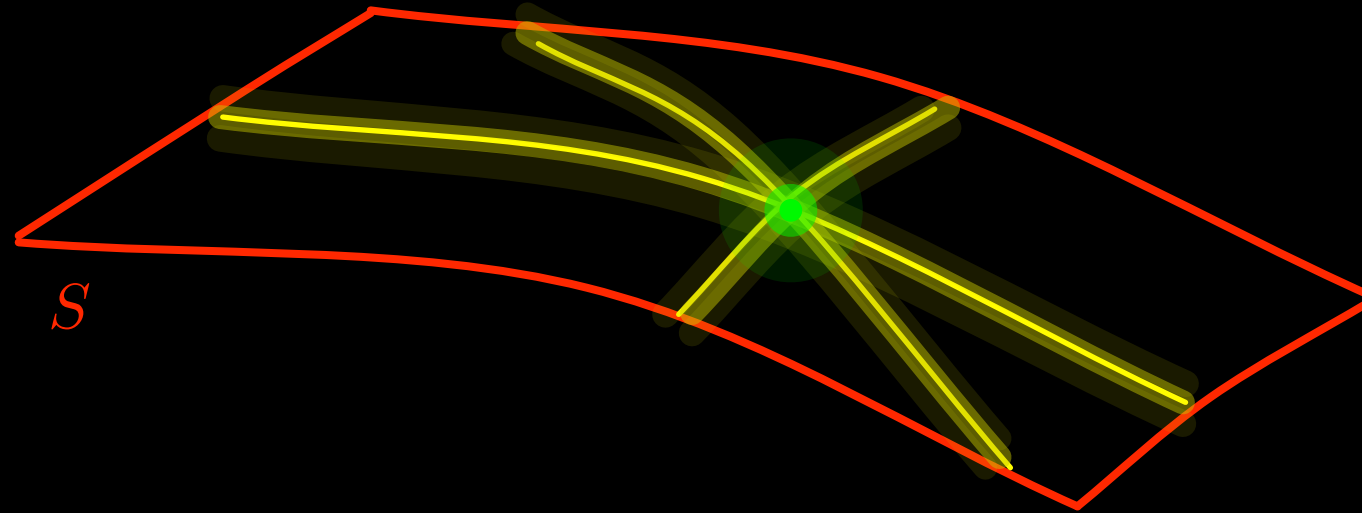


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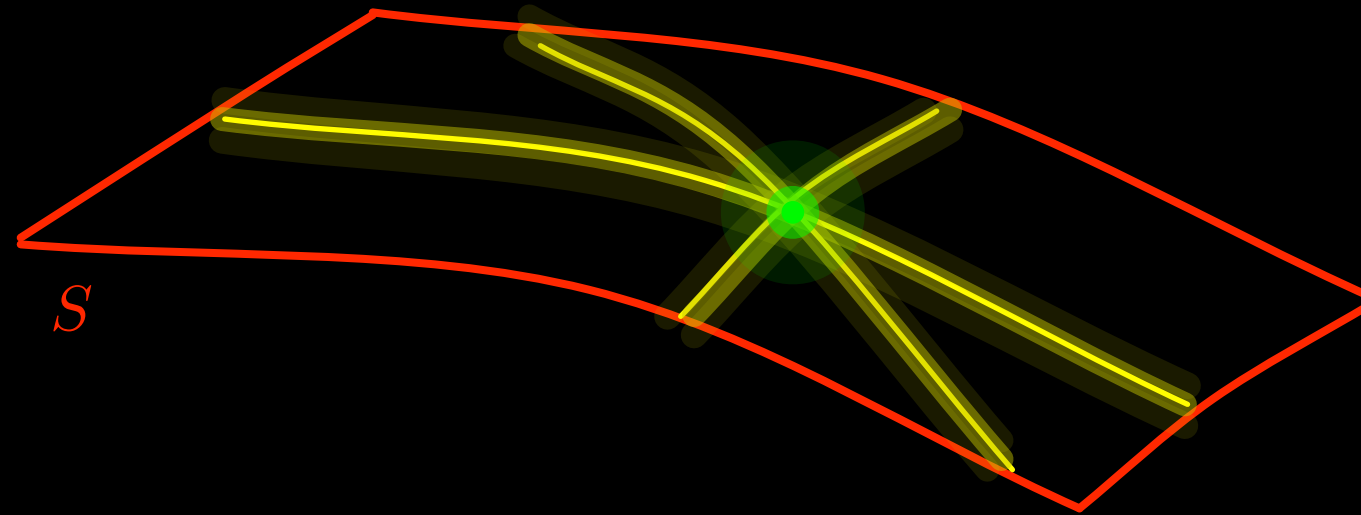


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 $\Rightarrow$  Higher Rank!

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Still topological: pairing of a quantum cohomology

# Yukawa as Quantum Residues

$$\theta = 0$$
$$W^{(ij)} = \oint dx \oint dy \frac{h_{12}^{(i)} h_{23}^{(j)}}{x y}$$

# Yukawa as Quantum Residues

$$\theta \neq 0$$

$$\begin{aligned}
 W^{(ij)} &= \oint dx \oint dy \frac{h_{12}^{(i)} h_{23}^{(j)}}{x y} + \mathcal{O}(\theta) \\
 &= \oint \oint \frac{h_{12}^{(i)}}{[\varphi^{(0)}, ]_{12,*}} \circledast \frac{h_{23}^{(j)}}{[\varphi^{(0)}, ]_{23,*}} \circledast (dx \wedge dy)
 \end{aligned}$$

$$\varphi_1^{(0)} \circledast \frac{h_{12}^{(i)}}{[\varphi^{(0)}, ]_{12,*}} - \frac{h_{12}^{(i)}}{[\varphi^{(0)}, ]_{12,*}} \circledast \varphi_2^{(0)} = h_{12}^{(i)} \quad , \quad \frac{h_{12}^{(i)}}{[\varphi^{(0)}, ]_{12,*}} = \frac{h_{12}^{(i)}}{x} + \mathcal{O}(\theta)$$

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**Maximal Rank, Hierarchy Given by Fluxes!**

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Natural way to get rank one+corrections. Might be relevant for F-theory phenomenology.

In particular, if N-C vanishes at intersecting point, numerology for CKM matrix very satisfactory.

But this should be taken with a grain of salt.

# Open Question

- Exceptional group?
- What structure of N-C is imposed on us?
- $SL(2, \mathbb{Z})$  monodromy
- Non-transversal intersection?